Unitary Equivalence of One-parameter Groups of Toeplitz and Composition Operators

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AMS Richmond, November 7, 2010

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Joint work with Eva Gallardo Gutiérrez, U. Zaragoza, Spain

The Hardy Hilbert space on the unit disk, $\mathbb{D} = \{z \in \mathbb{C} : |z| < 1\}$ is:

$$H^{2} = \{ f \text{ analytic in } \mathbb{D} : f(z) = \sum_{n=0}^{\infty} a_{n} z^{n} \text{ with } ||f||^{2} = \sum |a_{n}|^{2} < \infty \}$$

where for f and g in H^2 , we have $\langle f, g \rangle = \sum a_n \overline{b_n}$

and we will consider two types of operators on H^2 :

For ψ an analytic map of \mathbb{D} into the complex plane,

the analytic Toeplitz operator T_{ψ} is

$$(T_{\psi}f)(z) = \psi(z)f(z)$$
 for f in H^2

and, for φ an analytic map of \mathbb{D} into itself,

the composition operator C_{φ} is

$$(C_{\varphi}f)(z) = f(\varphi(z))$$
 for f in H^2

For example, if φ is defined by

$$\varphi(z) = \frac{3z+1}{z+3}$$

then φ is an automorphism of the disk \mathbb{D} with $\varphi(\pm 1) = \pm 1$ and $\varphi'(1) = \frac{1}{2}$

The spectrum of C_{φ} is the annulus

$$\sigma(C_{\varphi}) = \{\lambda : \frac{1}{\sqrt{2}} \le |\lambda| \le \sqrt{2}\}$$

and each λ with

$$\frac{1}{\sqrt{2}} < |\lambda| < \sqrt{2}$$

is an eigenvalue of infinite multiplicity for C_{φ} .

Similarly, if ψ is defined by

$$\psi(z) = \left(\frac{1-z}{1+z}\right)^{(i\log 2)/\pi}$$

then ψ is the covering map of the disk \mathbb{D} onto the annulus

$$\psi(\mathbb{D}) = \{\lambda : \frac{1}{\sqrt{2}} < |\lambda| < \sqrt{2}\}$$

The spectrum of the Toeplitz operator $T_{\overline{\psi}} = T_{\psi}^*$ is the annulus

$$\psi(\mathbb{D}) = \{\lambda : \frac{1}{\sqrt{2}} \le |\lambda| \le \sqrt{2}\}$$

and each λ with

$$\frac{1}{\sqrt{2}} < |\lambda| < \sqrt{2}$$

is an eigenvalue of infinite multiplicity for T_{ψ}^* .

Both of these operators are part of one-parameter groups of operators:

$$\varphi_t(z) = \frac{(1 + e^{-t})z + (1 - e^{-t})}{(1 - e^{-t})z + (1 + e^{-t})}$$

with
$$C_{\varphi_s}C_{\varphi_t} = C_{\varphi_s \circ \varphi_t} = C_{\varphi_{s+t}}$$
 for $-\infty < s, t < \infty$

and

$$\psi_t(z) = \left(\frac{1-z}{1+z}\right)^{it/\pi}$$

with
$$T_{\psi_s}^* T_{\psi_t}^* = T_{\psi_s \psi_t}^* = T_{\psi_{s+t}}^*$$
 for $-\infty < s, t < \infty$

and each λ with

$$e^{-t/2} < |\lambda| < e^{t/2}$$

is an eigenvalue of infinite multiplicity for each of C_{φ_t} and $T_{\psi_t}^*$.

IDEA!!

If there were a connection (e.g. similarity or unitary equivalence) between these operators,

then the eigenvectors for each of these operators should correspond to the eigenvectors for the same eigenvalue for the other operator!

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TRY infinitesimal generators!

The infinitesimal generator of a (semi)group A_t of operators

is an operator G such that for each f (in the domain of G)

$$Gf = \frac{d}{dt} \Big|_{t=0} A_t f$$

and analogous to the ideas from solution of first order linear elementary differential equations, we imagine that

$$A_t$$
 "=" e^{tG}

The infinitesimal generator of the group of composition operators is

$$\left(\frac{d}{dt}\Big|_{t=0} C_{\varphi_t} f\right)(z) = \frac{d}{dt}\Big|_{t=0} f(\varphi_t(z))$$

$$= f'(\varphi_t(z)) \frac{2e^{-t}(1-z^2)}{[(1-e^{-t})z + (1+e^{-t})]^2}\Big|_{t=0}$$

$$= f'(z) \frac{1-z^2}{2}$$

so G is the differential operator

$$(Gf)(z) = \frac{1}{2}(1 - z^2)f'(z)$$

A similar calculation gives the infinitesimal generator, H, of the group $T_{\psi_t}^*$.

Eigenvectors for the same eigenvalues of G and H should also be connected!

GOOD NEWS!!

The corresponding eigenspaces are 1-dimensional!

Let's try to match them up!

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For $-1/2 < \text{Re } \lambda < 1/2$, the eigenvectors of G and H are multiples of

$$w_{\lambda} = \left(\frac{1-z}{1+z}\right)^{-\lambda}$$
 and $v_{\lambda} = \left(1 - \frac{-i\sin\left(\lambda \frac{\pi}{2}\right)}{\cos\left(\lambda \frac{\pi}{2}\right)}z\right)^{-1}$

and, for each G and H,

the eigenvectors corresponding to $-1/2 < \lambda < 1/2$ have dense span in H^2

for
$$-1/2 < \lambda$$
, $\mu < 1/2$,

the relationship between w_{λ} and w_{μ} should be analogous to the relationship between v_{λ} and v_{μ} .

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BAD NEWS!!

Nasty computations give:

$$2\langle v_{\lambda}, v_{\mu} \rangle = \langle w_{\lambda}, w_{\mu} \rangle + 1$$

Not a good correspondence!

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AND it can't be fixed by multiplying the vectors by a constant!

If G and H are to correspond to each other, the relationship between w_{λ} and w_{μ} should be analogous to the relationship between v_{λ} and v_{μ} .

When does a unitary operator take a pair of one dimensional subspaces onto another (specific) pair??

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Lemma.

Let w_{λ} , w_{μ} , v_{λ} , and v_{μ} be non-zero vectors in \mathcal{H} , and let $M_{\lambda} = \operatorname{span}\{w_{\lambda}\}$, $M_{\mu} = \operatorname{span}\{w_{\mu}\}$, $N_{\lambda} = \operatorname{span}\{v_{\lambda}\}$, and $N_{\mu} = \operatorname{span}\{v_{\mu}\}$. There is a unitary operator U on \mathcal{H} such that $UM_{\lambda} = N_{\lambda}$ and $UM_{\mu} = N_{\mu}$ if and only if $\frac{|\langle w_{\lambda}, w_{\mu} \rangle|}{||w_{\lambda}|| ||w_{\mu}||} = \frac{|\langle v_{\lambda}, v_{\mu} \rangle|}{||v_{\lambda}|| ||v_{\mu}||}$

Lemma.

There is a unitary operator U on \mathcal{H} such that $UM_{\lambda} = N_{\lambda}$ and $UM_{\mu} = N_{\mu}$ if and only if

$$\frac{|\langle w_{\lambda}, w_{\mu} \rangle|}{\|w_{\lambda}\| \|w_{\mu}\|} = \frac{|\langle v_{\lambda}, v_{\mu} \rangle|}{\|v_{\lambda}\| \|v_{\mu}\|}$$

Theorem.

There is no unitary operator U on H^2 such that $U^*C_{\varphi_t}U = T_{\psi_t}^*$ for every real number t.

Proof:

Such a unitary U would take eigenspaces M_{λ} and M_{μ} onto the eigenspaces N_{λ} and N_{μ} , but for $\lambda = 0$ and $\mu = 1/4$, the angles don't match.

for
$$-1/2 < \lambda$$
, $\mu < 1/2$,

the relationship between w_{λ} and w_{μ} should be analogous to the relationship between v_{λ} and v_{μ} .

IDEA!!

Could the relationship

$$2\langle v_{\lambda}, v_{\mu} \rangle = \langle w_{\lambda}, w_{\mu} \rangle + 1$$

come from invariant subspaces??

Lemma.

For D bounded operator on Hilbert space \mathcal{H} and M an invariant subspace, then M^{\perp} is an invariant subspace for D^* .

$$D = \begin{pmatrix} A & B \\ 0 & C \end{pmatrix} \qquad D^* = \begin{pmatrix} A^* & 0 \\ B^* & C^* \end{pmatrix}$$

Lemma.

For D bounded operator on Hilbert space \mathcal{H} and M an invariant subspace, then M^{\perp} is an invariant subspace for D^* .

Furthermore, if r is an eigenvector for D with eigenvalue λ and r = p + q where p is in M and q is in M^{\perp} ,

then either q=0 or q is an eigenvector for the eigenvalue λ for the compression of D to M^{\perp} , which is the adjoint of the restriction of D^* to its invariant subspace M^{\perp} .

$$\begin{pmatrix} \lambda p \\ \lambda q \end{pmatrix} = \lambda \begin{pmatrix} p \\ q \end{pmatrix} = \lambda r = Dr = \begin{pmatrix} A & B \\ 0 & C \end{pmatrix} \begin{pmatrix} p \\ q \end{pmatrix} = \begin{pmatrix} Ap + Bq \\ Cq \end{pmatrix}$$

so
$$\lambda q = Cq$$

Let [1] denote subspace of H^2 spanned by constants so that $H^2 = [1] \oplus zH^2$.

For each t, the subspace zH^2 is an invariant subspace for T_{ψ_t} and for $C_{\varphi_t}^*$.

Letting $x_{\lambda} = w_{\lambda} - 1$ and $u_{\lambda} = v_{\lambda} - 1$, each in zH^2 , work above means that x_{λ} and u_{λ} are each eigenvectors of the compressions of C_{φ_t} and $T_{\psi_t}^*$ to zH^2 and they are eigenvectors of the compressions of G and H to zH^2 that correspond to the same eigenvalues.

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 zH^2 is an invariant subspace for T_{ψ_t} and for $C_{\varphi_t}^*$ for each t.

Letting $x_{\lambda} = w_{\lambda} - 1$ and $u_{\lambda} = v_{\lambda} - 1$, each in zH^2 , this means that x_{λ} and u_{λ} are each eigenvectors of the compressions of C_{φ_t} and $T_{\psi_t}^*$ to zH^2 and they are eigenvectors of the compressions of G and H to zH^2 that correspond to the same eigenvalues.

A MIRACLE:

$$2\langle u_{\lambda}, u_{\mu} \rangle = \langle x_{\lambda}, x_{\mu} \rangle$$

Theorem.

(1) The operator U defined by

$$U(x_{\lambda}) = \sqrt{2}u_{\lambda}$$

can be extended to a unitary operator of zH^2 onto itself.

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can be extended to a unitary operator of zH^2 onto itself.

(2) For each real number t,

$$U \left. C_{\varphi_t}^* \right|_{zH^2} = \left. T_{\psi_t} \right|_{zH^2} U$$

so the operators $C_{\varphi_t}^*|_{zH^2}$ and $T_{\psi_t}|_{zH^2}$ are unitarily equivalent.

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(2) For each real number t,

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so the operators $C_{\varphi_t}^*|_{zH^2}$ and $T_{\psi_t}|_{zH^2}$ are unitarily equivalent.

That is, there is a unitary operator on zH^2 that shows the restrictions of $C_{\varphi_t}^*$ and T_{ψ_t} to zH^2 are unitarily equivalent for each t.

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Slides posted on webpage:

www.math.iupui.edu/~ccowen/Richmond1011.pdf

Paper posted on webpage:

www.math.iupui.edu/~ccowen/Downloads.html